

References

- ALLPRESS, J. G. (1965). *J. Inorg. Nucl. Chem.* **27**, 1521.
 BACHMANN, H. G. & ZEMANN, J. (1961). *Acta Cryst.* **14**, 474.
 BYSTRÖM, A. & WESTGREN, A. (1943). *Ark. Kemi Min. Geol.* **16B**, no. 14.
 CROMER, D. T., LARSON, A. C. & ROOF, R. B. (1964). *Acta Cryst.* **17**, 272.
 CRUICKSHANK, D. W. J. (1960). *Acta Cryst.* **13**, 774.
 DAUBEN, C. H. & TEMPLETON, D. H. (1955). *Acta Cryst.* **8**, 841.
 FRONDEL, C. & BARNES, I. (1958). *Acta Cryst.* **11**, 562.
 HOWELLS, E. R., PHILLIPS, D. C. & ROGERS, D. (1950). *Acta Cryst.* **3**, 210.
International Tables for X-ray Crystallography (1962). Vol. III, pp. 202, 206, 212. Birmingham: Kynoch Press.
 KAY, M. I. (1961). *Acta Cryst.* **14**, 80.
 KLUG, H. P. & ALEXANDER, L. E. (1954). *X-ray Diffraction Procedures*, p. 376. New York: John Wiley.
 KOVBA, L. M., POLUNINA, G. P., SIMANOV, YU. P. & IPOLITOVA, B. A. (1961). *Issledovania v Oblasti Khimii Urana*, p. 15. Moscow Univ. Press.
 LECIEJEWICZ, J. (1961). *Acta Cryst.* **14**, 66.
 LIPSON, H. & WOOLFSON, M. M. (1952). *Acta Cryst.* **5**, 680.
 MILLS, O. S. & ROLLETT, J. S. (1961). *Computing Methods and the Phase Problem in X-ray Crystal Analysis*, p. 117. Oxford: Pergamon Press.
 RIETVELD, H. M. (1966). *Acta Cryst.* **20**, 508.
 ROOF, R. B. (1961). *Acta Cryst.* **14**, 934.
 ROOF, R. B., CROMER, D. T. & LARSON, A. C. (1964). *Acta Cryst.* **17**, 701.
 RÜDORFF, W. & PFITZER, F. (1954). *Z. Naturforsch.* **96**, 568.
 SAMSON, S. & SILLÉN, L. G. (1947). *Ark. Kemi Min. Geol.* **25A**, no. 21.
 SLEIGHT, A. W. & WARD, R. (1962). *Inorg. Chem.* **1**, 790.
 ZACHARIASEN, W. H. (1954). *Acta Cryst.* **7**, 795.

Acta Cryst. (1967). **23**, 272

The Crystal and Molecular Structure of *meso*-3,3'-Di-(*p*-bromophenyl)bi-3-phthalidyl

BY V. KALYANI, H. MANOHAR* AND N. V. MANI†

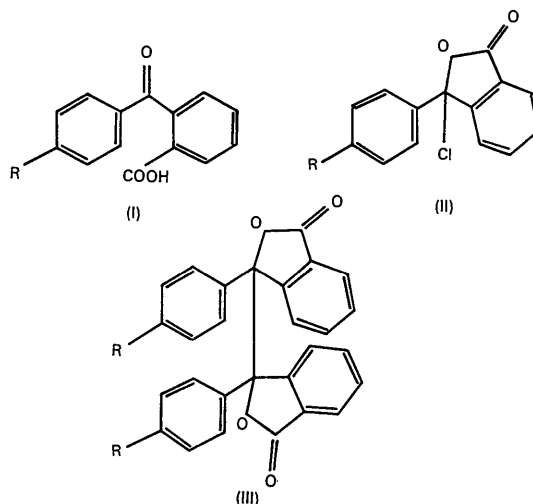
Department of Physics, Indian Institute of Science, Bangalore 12, India

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The crystal structure of the title compound has been determined by the heavy-atom method and Fourier techniques using visually estimated photographic data. It crystallizes in the space group $P\bar{1}$ with cell dimensions $a=7.98$, $b=8.08$, $c=9.66$ Å, $\alpha=85^\circ 23'$, $\beta=98^\circ 24'$ and $\gamma=104^\circ 41'$. The parameters have been refined by three-dimensional least-squares procedures with anisotropic thermal vibrations for all the atoms. The final R index for 932 observed reflexions is 0.096. It has been established that the stereoisomer taken up for investigation is the *meso* form existing in a staggered configuration. Two characteristics of the lactone group, *viz.*, the inequality of the C–O bonds and the planarity of the group are also confirmed in the present study.

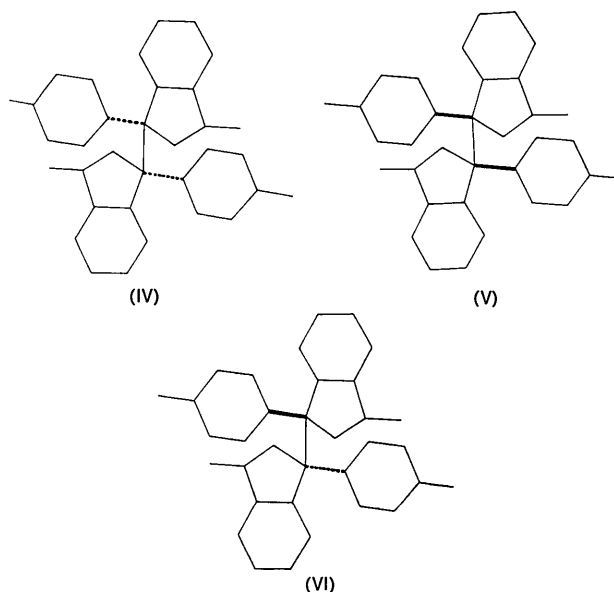
Introduction

A variety of reducing agents convert *o*-benzoylbenzoic acids (I) and their acid chlorides (II) to 3,3'-diarylbi-phthalidyls (III). Two isomers, *viz.* the racemic (DL) [IV and V, one being the *dextro* (D) and the other the *laevo* (L) form] and the *meso* form (VI) having structure (III) are theoretically possible. However, no attempt has so far been made to define their stereochemistry. This information would be of considerable value in understanding the steric factors involved in their formation, particularly because the two stereoisomers are formed in unequal proportions. In view of this, the more



* Present address: Department of Inorganic and Physical Chemistry, Indian Institute of Science, Bangalore 12, India.

† Present address: Laboratorium voor Structuurchemie, Rijksuniversiteit Groningen, Bloemensingel 10, Groningen, The Netherlands.



abundant isomer of the bromo derivative (III, 4'R = Br) was taken up for X-ray analysis, the results of which are reported below.

Crystal data

The halo derivatives and the unsubstituted diphenylbipthalidyl (III, 4'R = Br, Cl and H) were prepared by the reaction of *o*-benzoylbenzoyl chlorides (II, 4'R = Br, Cl and H) with a solution of sodium iodide in acetone. The stereoisomers, *meso* and racemic, were separated by fractional crystallization. The more abundant fractions of the three compounds were kindly made available to us by Prof. M. V. Bhat of the Organic Chemistry Department. The crystals were transparent and most of the specimens were tabular or needle-shaped. The crystal data for all the three compounds are summarized in Table 1. The cell dimensions were determined from rotation and zero-level Weissenberg

photographs taken with Cu $K\alpha$ radiation for crystals rotating about the three crystallographic axes. Density determinations were made by flotation in aqueous potassium iodide solution.

The good agreement in axial dimensions and inter-axial angles (Table 1) for the triclinic crystals of bromo and chloro compounds as well as similarity in their crystal habits suggested a high degree of isomorphism between the two compounds. Statistical tests (Wilson, 1949; Howells, Phillips & Rogers, 1950) carried out with corrected intensities in $h0l$ and $0kl$ zones to distinguish between the space groups $P1$ and $P\bar{1}$ were ambiguous. This is probably due to the presence of heavy atoms and also to the inadequate number of observations used to draw the $N(z)$ distribution curves.

The bromine substituted title compound (III, 4'R = Br) was first taken up for X-ray analysis and the heavy atom method was employed for the solution of the structure.

Data collection

The three zero-levels were recorded with Cu $K\alpha$ radiation by the multiple-film Weissenberg technique. Higher level integrated intensities of the general hkL reflexions, for L from 1 through 8, were collected by the equi-inclination method. Of the 2720 possible independent reflexions accessible in the effective copper sphere, 1913 were recorded, out of which 932 were in the measurable range. The intensities were estimated visually by means of a time-exposure calibrated strip prepared with the same specimen and corrected for Lorentz and polarization factors. The physical corrections for absorption for cylindrical specimens of mean radii 0.20, 0.22 and 0.32 mm were applied to the $0kl$, $h0l$ and hkl data respectively (Bond, 1959). Wilson's (1942) procedures were used to derive the absolute scale for all the three zero-level data and comparison of common reflexions yielded the scale factors to be used for the conversion of the higher level data.

Table 1. Crystal data for diphenylbipthalidyl and its halo derivatives

	Unsubstituted, $C_{28}H_{18}O_4$	Bromo derivative, $C_{28}H_{16}O_4Br_2$	Chloro derivative, $C_{28}H_{16}O_4Cl_2$
Formula weight	418.42	576.24	487.32
Cell parameters			
a	7.72	7.98	7.89
b	14.38	8.08	8.06
c	9.93	9.66	9.65
α	—	85.4	88.7
β	98.5 \pm 0.9°	98.4	99.5
γ	—	104.7	104.4
Unit cell volume, U	1090.4 \AA^3	595.2 \AA^3	585.8 \AA^3
Space group	$P2_1/c$ (no. 14)	$P\bar{1}$ (no. 2)	$P\bar{1}$ (no. 2)
Density			
D_m	1.288 g.cm $^{-3}$	1.619 g.cm $^{-3}$	1.371 g.cm $^{-3}$
D_c	1.274	1.607	1.381
Molecules per unit cell,			
Z	2	1	1
$F(000)$	436 e	286 e	250 e
Linear absorption coefficient for X-rays ($\lambda = 1.542 \text{\AA}$)	6.98 cm $^{-1}$	49.99 cm $^{-1}$	26.96 cm $^{-1}$

Structure determination

The analysis of the structure was commenced assuming the centrosymmetric space group. With one formula unit in the full cell, the space group requirement will be fulfilled only if the molecule has a centre of inversion as its symmetry element. The easy convergence of the refinement of the structure in this space group to a final R value of 0.096 confirmed the validity of our assumption and also definitely established that the substance under investigation was indeed a *meso* isomer. Further support is received from the fact that the abundant fraction of the free diphenylbiphtalidyl crystallizes in the uniquely derivable space group $P2_1/c$. With only two formula units of this compound in the monoclinic cell and a fourfold multiplicity of a general position, considerations of space group symmetry again require that the molecule should possess a centre of symmetry. Therefore the assumption of the same molecular symmetry for the abundant isomer of the bromo compound was amply justified.

The bromine positions were easily found by interpretation of (010) and (001) Patterson projections made using 84 and 75 observed reflexions respectively. The Fourier syntheses based on the bromine phases constructed subsequently with 78 and 71 amplitudes from $h0l$ and $hk0$ zones confirmed the heavy atom positions. The first set of structure factor calculations yielded the R values 0.24 and 0.31 for the two respective zones, when 16 and 13 lighter atoms, placed on the peaks, were included in the calculations.

Successive approximations to the electron densities in these projections initiated with phases given by the bromine and refined by continually adding lighter atoms to the trial model gave a reasonably accurate picture of the molecular structure. The trial structure at this stage included all the atoms other than hydrogen atoms and the atomic parameters were refined by an iteration process of residual electron density projections down the b and c axes. The refinement using zonal data was terminated when the R values did not drop below 0.22 and 0.19.

For the foregoing calculations, the atom form factors were derived with the use of an analytical expression involving five constants for each atom (Forsyth & Wells, 1959) and a uniform isotropic B value of 3.5 \AA^2 was applied to all atoms. The above mentioned Fourier summations at 6° intervals were done on the Ferranti SIRIUS computer at N.A.L., Bangalore, with programs devised in this laboratory (Kannan, Vijayan & Nambudiri, 1966).

Three-dimensional refinement of the structure by the least-squares method was then attempted with all 932 observed reflexions. The initial stage of the refinement, which involved the diagonal approximation and use of individual isotropic thermal parameters, was done on the CDC-3600 computer system at Bombay with a program prepared by Kannan *et al.* (1966). The atom form factors were now modified by use of a nine-

parameter analytical expression evolved by Cromer & Waber (1965). Three separate sets, each consisting of three cycles of structure-factor least-squares (SFLS) calculations, lowered the overall R value from 0.30 to 0.185. The unusual feature in these calculations was the consistently low values ranging from 1.5 to 2.5 \AA^2 of the isotropic B values indicated for all atoms. These were increased at the end of every cycle to values between 3.0 and 4.5 \AA^2 and the various zones were individually scaled up before another run was made. The low values of the thermal parameters encountered here may be attributed to appreciable variations in absorption factors from the mean due to the large absorption coefficient of 50 cm^{-1} ($\lambda = 1.542 \text{ \AA}$) and consequently, the final temperature factors reported are probably lower than the true ones. Bond distances and valency angles computed for the first time at this stage revealed considerable distortions in the phenyl ring with C—C lengths varying from 1.2 to 1.6 \AA and internal angles from 108 to 130° . Atomic positions of the carbons C(1) to C(6) (Fig. 1) were then adjusted by trial to bring the distances in the phenyl ring to normal values.

The last set of structure factor calculations showed some consistent discrepancies between F_o and F_c values which suggested introduction of anisotropic (ellipsoidal) thermal parameters for some of the atoms. Indeed difference electron density diagrams constructed with $h0l$ and $hk0$ data earlier clearly indicated an anisotropic vibrational pattern for the bromine atom. The final cycles of SFLS refinement were performed on the National Elliott 803-B computer at Hindustan Aeronautics Ltd, Bangalore, with a generalized program devised by Dr G.A. Mair of the Royal Institution, London. This program minimizes the function $\sum w|kF_o - F_c|^2$, where w is the weight assigned to the observed structure factors and k is a refinable scale correction. This program uses the block-diagonal approximation with a 3×3 matrix for the positional parameters and a 6×6 or 1×1 matrix for the anisotropic or isotropic thermal parameters of each atom. The scale factor for the F_o 's and the overall isotropic B

Table 2. Fractional atomic coordinates
e.s.d.'s $\times 10^3 \text{ \AA}$ in parenthesis)

	<i>x</i>	<i>y</i>	<i>z</i>
Br	0.5373 (2)	0.2429 (2)	0.5666 (2)
O(1)	-0.3576 (11)	0.2592 (11)	-0.0204 (12)
O(2)	-0.1562 (10)	0.1097 (10)	0.0542 (10)
C(1)	0.3739 (18)	0.1847 (17)	0.4046 (16)
C(2)	0.4429 (17)	0.1971 (16)	0.2750 (14)
C(3)	0.3242 (15)	0.1712 (15)	0.1539 (16)
C(4)	0.1425 (14)	0.1270 (15)	0.1584 (14)
C(5)	0.0753 (17)	0.1161 (15)	0.2883 (15)
C(6)	0.1970 (18)	0.1451 (17)	0.4126 (16)
C(7)	0.0132 (13)	0.0939 (14)	0.0232 (14)
C(8)	0.0531 (14)	0.2311 (14)	-0.0952 (13)
C(9)	0.2005 (16)	0.2747 (16)	-0.1720 (14)
C(10)	0.1992 (18)	0.4068 (16)	-0.2747 (16)
C(11)	0.0626 (17)	0.4901 (16)	-0.3026 (15)
C(12)	-0.0817 (18)	0.4450 (16)	-0.2251 (15)
C(13)	-0.0827 (14)	0.3100 (15)	-0.1235 (13)
C(14)	-0.2150 (18)	0.2328 (17)	-0.0271 (16)

ring and finally extended to all atoms. Following seven cycles of SFLS calculations, the refinement was terminated when the R value dropped to 0.096 from an initial value of 0.185 for all the observed reflexions. Table 3 lists the F_o and F_c values based on the final atomic positions given in Table 2. Structure factors have been computed for about 980 unobserved reflexions (not given), of which only 60 are slightly above the threshold amplitudes.

A disconcerting feature of the refinement procedures which persisted till the end concerns the large parametral shifts shown for some of the light atoms particularly those of the phenyl ring. Considerable fudging of the shifts had to be exercised in these cases for every

cycle, as, otherwise, the indicated movements would have resulted in a highly distorted geometry of the phenyl ring. The reason for this is not quite clear, although it is believed to be partly due to the use of 'truncated' three-dimensional data and partially diagonalized matrix elements. A final three-dimensional difference map (not shown) was computed in an attempt to locate the hydrogen atoms; but it did not prove useful in indicating their positions. Hence the hydrogen atoms were located at a calculated distance of 1.08 Å from the appropriate carbons. Addition of all the hydrogen atoms to the model lowered R by only 0.1%. A temperature factor of 4.5 Å² was used for all the H atoms; neither this nor their positional parameters was refined. The mean parametral shift in the light atom position in the ultimate least-squares cycle was 0.021 Å ($\sim 1.4\sigma$), the only exceptions being the x coordinate of C(1) and the y coordinate of C(6) for which the shifts measured as much as 2.6σ . The shifts in the heavy atom position was 0.0006 Å ($\sim 0.3\sigma$). The e.s.d.'s (σ 's) are the average values obtained from the inverse of the appropriate block-matrix A_{mm}^{-1} and from the relation

$$\sigma_{mm} = [A_{mm}^{-1} \sum w(|kF_o - F_c|)^2 / (n-s)]^{1/2},$$

where the symbols have the usual meanings.

Results

The final atomic coordinates (in fractional units) that resulted from our analysis and the e.s.d.'s (in Å) (in parenthesis) are given in Table 2. The components b_{ij} of the anisotropic thermal parameters and the equivalent isotropic B values (Å²) as evaluated from Hamilton's (1959) formula

$$B = \frac{4}{3} \sum_i \sum_j b_{ij}(\mathbf{a}_i \cdot \mathbf{a}_j)$$

are represented in Table 4. The magnitudes and direction cosines of the principal axes of thermal vibration

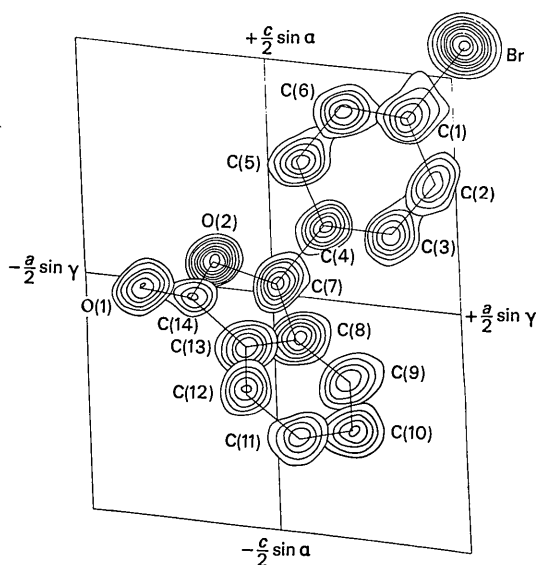


Fig. 1. Sections of the final three-dimensional electron density map projected on the (010) plane. As the molecule possesses an inversion centre, only one half is represented. Contours start at 1.0, 1.5 and 3.0 e.Å⁻² for C, O and Br respectively. Intervals are at 0.5 e.Å⁻³ for C and O and 2 e.Å⁻³ for Br.

Table 4. Anisotropic thermal parameters ($\times 10^4$) for the expression

$$T = \exp[-(b_{11}h^2 + b_{22}k^2 + b_{33}l^2 + b_{23}kl + b_{31}hl + b_{12}hk)]$$

and the equivalent isotropic thermal parameters (Hamilton, 1959)

Atom	b_{11}	b_{22}	b_{33}	b_{12}	b_{23}	b_{13}	B
Br	256	272	119	184	-79	-113	5.73 Å ²
O(1)	120	219	163	159	101	51	4.61
O(2)	105	108	123	44	-1	122	3.08
C(1)	269	139	109	80	-19	-75	4.77
C(2)	178	181	64	118	-20	-33	3.66
C(3)	112	142	118	-19	-55	73	3.54
C(4)	105	126	69	54	28	14	2.73
C(5)	190	126	74	115	32	85	3.28
C(6)	207	185	97	49	-30	118	4.27
C(7)	67	127	73	73	20	43	2.38
C(8)	105	99	61	32	37	72	2.37
C(9)	169	174	56	12	38	137	3.45
C(10)	228	139	94	95	29	69	4.03
C(11)	192	173	72	73	-13	48	3.79
C(12)	245	131	77	71	-27	8	3.99
C(13)	119	142	52	18	-13	80	2.72
C(14)	190	180	117	-93	-59	70	4.68

ellipsoids, as derived from the parameters of Table 4, are listed in Table 5. Although serious attempts to interpret these values are not made in view of their limited accuracy mentioned earlier, the maximum motion appears to be in a direction normal to the 'phthal-

idyl' ring. The final three-dimensional electron density map corresponding to the asymmetric half of the molecule is shown in Fig. 1 by means of superimposed contour sections drawn perpendicular to *b* axis. The atomic arrangement is also indicated.

Discussion

Table 5. *Magnitudes and direction cosines, relative to the reciprocal crystallographic axes, of the principal axes of the thermal vibration ellipsoids*

Atom	Axis <i>i</i>	B_i (Å ²)	g_{i1}	g_{i2}	g_{i3}
Br	1	2.98	0.7769	0.1593	0.7257
	2	8.56	-0.5984	0.1158	0.6459
	3	6.00	0.1977	0.9813	-0.2373
O(1)	1	7.49	0.7771	0.7630	0.2702
	2	2.22	-0.2310	0.5156	-0.8422
	3	4.32	-0.5856	0.3902	0.4679
O(2)	1	4.93	0.8690	0.3017	0.5879
	2	1.88	0.4454	0.3897	-0.8028
	3	2.36	-0.2148	0.8730	0.1001
C(1)	1	7.34	0.5686	0.4451	-0.7028
	2	3.21	0.5553	0.6556	0.6590
	3	3.36	-0.5925	0.6278	-0.2523
C(2)	1	1.95	-0.6740	0.2876	-0.6907
	2	5.52	-0.2121	0.8298	0.2538
	3	4.13	0.5613	0.5385	-0.6689
C(3)	1	4.80	-0.5318	0.6291	-0.4493
	2	2.26	0.6416	0.2795	-0.6751
	3	3.47	0.5536	0.7261	0.5855
C(4)	1	3.47	0.6971	0.8289	-0.1898
	2	1.85	-0.6909	0.5308	-0.1158
	3	2.46	0.1875	0.1781	0.9754
C(5)	1	4.43	0.5398	0.4128	0.8296
	2	2.42	-0.8329	-0.1631	0.4342
	3	2.89	-0.1255	0.8965	-0.3536
C(6)	1	1.72	0.8824	0.3942	-0.3282
	2	6.47	-0.4332	0.5462	-0.7025
	3	4.88	0.1832	0.7405	0.6312
C(7)	1	3.36	0.2780	0.8954	0.4050
	2	1.22	-0.1595	0.4400	-0.9050
	3	2.05	-0.9473	-0.0705	0.1335
C(8)	1	1.90	-0.8774	0.0584	0.2465
	2	3.15	0.4508	0.3131	0.9005
	3	2.70	0.1643	0.9484	-0.3596
C(9)	1	1.72	-0.8756	0.0223	0.2788
	2	4.33	0.4613	0.7744	0.5672
	3	3.70	-0.1454	0.6324	-0.7755
C(10)	1	2.56	-0.7994	0.0219	0.4293
	2	5.35	-0.4304	0.3149	-0.8733
	3	3.96	0.4197	0.9492	0.2324
C(11)	1	5.18	-0.0793	0.8357	-0.5397
	2	2.85	0.9854	0.3781	0.2136
	3	3.85	-0.1504	0.3996	0.8152
C(12)	1	5.86	-0.2027	0.0096	0.9352
	2	2.42	0.7575	0.8085	0.1640
	3	3.88	-0.6205	0.5889	-0.3146
C(13)	1	1.43	0.8402	0.3641	-0.4153
	2	3.66	-0.0039	0.9259	0.2143
	3	2.97	-0.5428	0.1046	-0.8846
C(14)	1	6.82	-0.4070	0.6329	-0.6205
	2	3.02	-0.1127	0.4726	0.7839
	3	3.81	0.9065	0.6137	-0.0367

Interatomic bond lengths and valency angles in the molecule with the corresponding e.s.d.'s in parenthesis are indicated in Fig. 2. The mean value of the aromatic C-C bond is 1.406 Å, which is slightly above the normal value of 1.395 Å. However, since the standard deviation of this bond is 0.022 Å, by Cruickshank's significance criterion (Lipson & Cochran, 1957), the difference is not appreciable. None of the bond lengths between two sp^3 hybridized carbon atoms (1.571 Å), between sp^3 and sp^2 (1.538 Å) and between two sp^2 hybridized carbon atoms (1.494 Å) differ significantly from accepted values of 1.545, 1.525 (Sutton, 1958) and 1.477 Å (Dewar & Schmeising, 1959) respectively. The C(aromatic)-Br distance of 1.893 Å is normal and agrees well with the average of several recent accurate determinations.

The two nominally equivalent C-O bonds differ significantly in length. Similar differences were first encountered in the γ -lactone ring in himbacine hydrobromide (Fridrichsons & Mathieson, 1962) and have subsequently been authenticated from many other determinations (Karle & Karle, 1966; Gabe, 1962; Schoenborn & McConnell, 1962; Fridrichsons, Mathieson & Sutor, 1963; Kartha & Haas, 1964; Asher & Sim, 1965). The C(7)-O(2) bond length of 1.465 Å does not differ significantly from the C-O single bond distance of 1.43 Å (Sutton, 1958). On the other hand, the C(14)-O(2) bond adjacent to the C=O group is considerably shorter, being 1.353 Å. This shortening has been interpreted in two ways (McConnell, Mathieson & Schoenborn, 1964): either as due to the hybridization of O(2) which is assumed to be in a trigonal sp^2 state and its interaction with C(14); or, alternatively, as due to a partial double bond character of the C(14)-O(2) bond, with O(2) acquiring a positive charge and O(1) a negative charge. This accounts for the planarity of the atoms C(13), C(14), O(1), O(2) and C(7) of the lactone group, just as in the case of the ester group and carboxylic acids. These two characteristics, inequality of the two C-O bonds and planarity of the lactone group, have been confirmed in the present studies. The amount of double bond character of the C(14)-O(2) bond has been estimated from a plot of double bond character *versus* C-O bond lengths, as indicated by Curl (1959) and is found to be approximately 14%. In common with other structures, no evidence for the slight lengthening of the C(14)=O(1) double bond, as suggested by Curl, has been noticed, the bond distance of 1.221 Å ($\sigma=0.021$ Å) agreeing well with the normal value of 1.23 Å (Sutton, 1958) within experimental limits.

The average value of the C–C–C angle in the two phenyl rings is 119.9° . The following points regarding the angles around the sp^2 hybridized carbon atom C(14) may be observed. Although the sum of the angles at C(14) is 360° , the two exocyclic angles are appreciably different from the internal angle of 107.5° . The internal angle in the lactone ring, *i.e.* not containing the double bond, is the smallest. Further O(1)–C(14)–O(2), the angle subtended by the two oxygen atoms, is smaller than O(1)–C(14)–C(13) by 8.6° . The latter interesting feature has been observed by us in several of the structures containing the γ -lactone ring system. A few examples of the differences between the two angles observed in various structures are as follows: (i) annotinine bromohydrin (Przybylska & Ahmed, 1958), 6° ; (ii) himbacine hydrobromide (Fridrichsons & Mathieson, 1962), 7° ; (iii) bromodilactone from jacobine (Mathieson & Taylor, 1963), 8° and 13° ; (iv) *p*-bromobenzoate-glauucarubin (Kartha & Haas, 1964), 7° ; (v) shellolic bromolactone hydrate (Gabe, 1962), 6° and (vi) anemonin (Karle & Karle, 1966), 5° and 8° . There are some significant differences between the tetrahedral angles at C(7). The widening of the angle C(4)–C(7)–C(8) is the result of steric hindrance between the hydrogen atoms attached to atoms C(3) and C(9), this close

contact being 2.22 \AA . The decrease in the angle O(2)–C(7)–C(8) from 109.5° to 102.5° may be ascribed to the considerable strain in the γ -lactone ring as a result of the shortening of one C–O bond.

Since the bromine atom and the carbon atom C(1) to which it is attached have high temperature factors, it was felt that thermal motion might have considerable effect on the C–Br bond length. This distance was therefore corrected for thermal oscillations using the method of Busing & Levy (1964). The motions of C and Br cannot be described completely as either correlated or uncorrelated. However, the former description may be a better approximation to physical reality, as the C–Br bond is a strong covalent linkage. Therefore, the correction was applied assuming correlated or 'in phase' motion. It was however found that the corrected distance of 1.903 \AA is not much different from the uncorrected value. The C(14)–O(1) distance was also corrected likewise, the corrected value being 1.225 \AA . Here again the correction is not significant, being much smaller than the e.s.d. of the bond.

The equations for the mean planes through the phenyl ring, the fused ring system and the lactone group have been derived by Blow's method (Blow, 1960). The equations are referred to a set of orthogonal

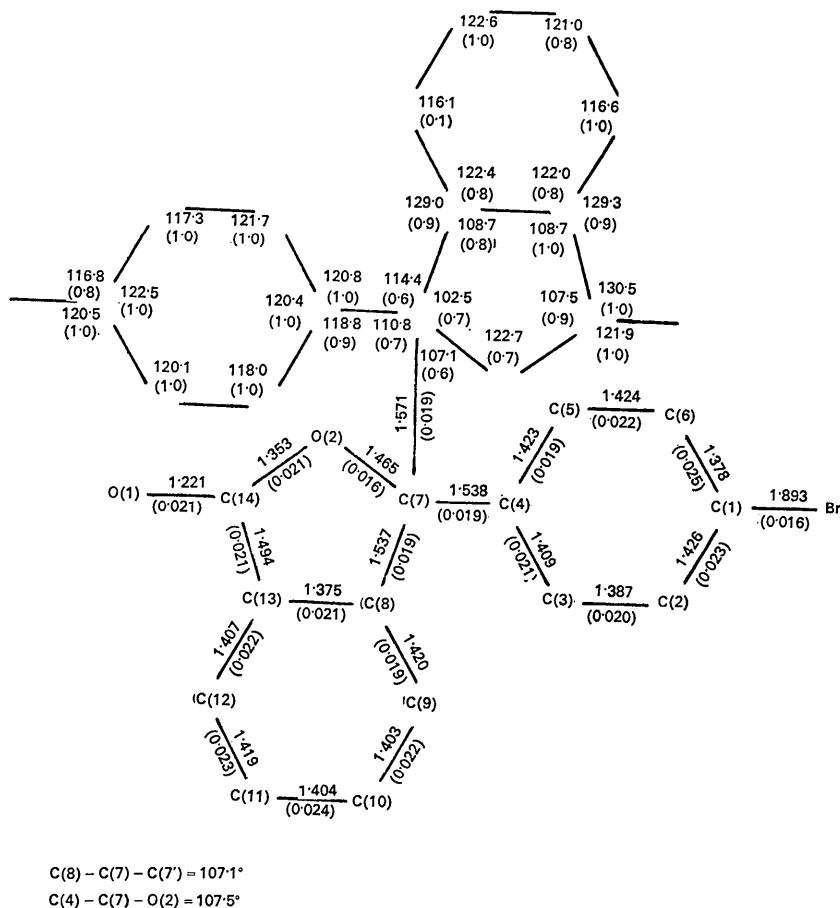


Fig. 2. Bond lengths (Å) and bond angles (°); values in parenthesis denote the estimated standard deviations.

axes X' , Y' , Z' where X' coincides with the crystallographic X direction, Y' lies in the XY plane and Z' is perpendicular to the XY plane. The equations are as follows:

Phenyl ring:

$$0.2303X' - 0.9678Y' + 0.1017Z' + 0.7201 = 0$$

Fused ring system:

$$0.2180X' + 0.7041Y' + 0.6758Z' - 0.6376 = 0$$

Lactone group:

$$0.1937X' + 0.7097Y' + 0.6773Z' - 0.6673 = 0$$

A projection of the molecule (excluding hydrogen atoms) in the plane of the fused ring system and the out-of-plane distances of the atoms are given in Fig. 3. The phenyl ring, the fused ring system and the lactone group are essentially planar, the displacement of the atoms from the respective mean planes being, in general, less than their standard deviations normal to the planes. However, the bromine and the oxygen atom O(1) are off the mean plane of the phenyl ring and the fused ring system by 0.116 and 0.080 Å respectively.

The phenyl ring is inclined at an angle of 124.2° to the plane of the fused ring system belonging to the same half of the molecule. This orientation of the benzene plane and the centrosymmetrical conformation of the molecule correspond to a position of least steric hindrance. Any deviation from this configuration would only shorten some of the intramolecular contacts while increasing the others. The short intramolecular non-bonded distances are given in Table 6.

Table 6. *Short intermolecular contacts*

BrA	—C(1')D	3.681 Å
BrA	—C(10)C	3.838
C(2)A	—C(10')B	3.709
C(2)A	—O(1)B	3.409
C(3)A	—O(1)B	3.151
C(1')A	—C(9')B	3.500
C(12)A	—C(4')E	3.729
C(12)A	—C(5')E	3.661
C(11)A	—C(5')E	3.638

The coordinates of the atoms of molecule *A* are related to those of molecules *B* to *E* by the following symmetry relations:

Molecule	Coordinates
<i>A</i>	$x \quad y \quad z$
<i>B</i>	$1+x \quad y \quad z$
<i>C</i>	$x \quad y \quad 1+z$
<i>D</i>	$1+x \quad y \quad 1+z$
<i>E</i>	$x \quad 1+y \quad z$

The arrangement of molecules in the unit cell as seen in the $h0l$ and $hk0$ projections is shown in Figs. 4 and 5. It can be seen that the structure is based purely on packing considerations which are decided by the shape of the molecule, fragments of one filling the concavity of its neighbours. The close intermolecular contacts less than 4 Å are marked in the two figures and also listed in Table 7. None of them is abnormal.

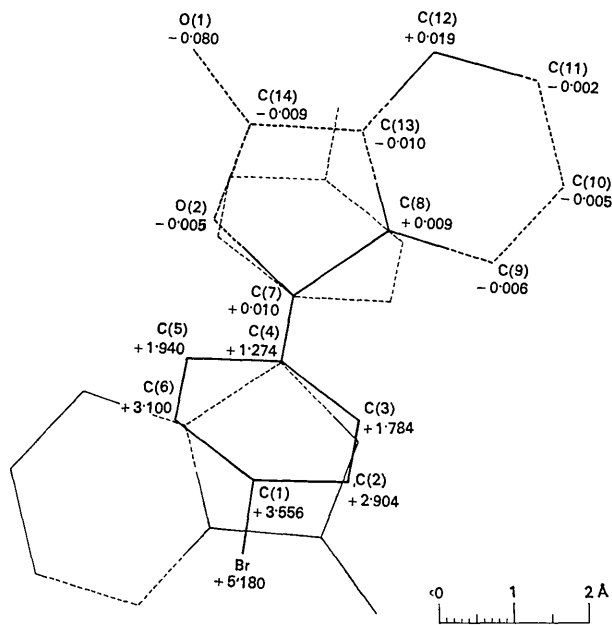


Fig. 3. Mean plane of the 'phthalidyl' ring system showing the perpendicular distances (Å) of the atoms of one half molecule from the plane.

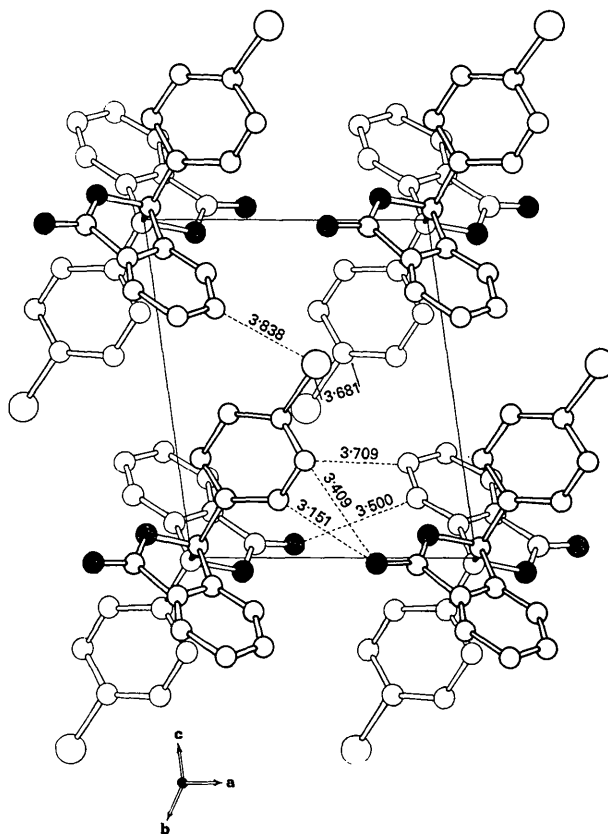


Fig. 4. View of the structure along $[010]$ showing the packing arrangement and short intermolecular contacts. The hydrogen atoms are not shown. The large open circles denote the bromine atoms and the smaller open and black circles represent the carbon and oxygen atoms respectively.

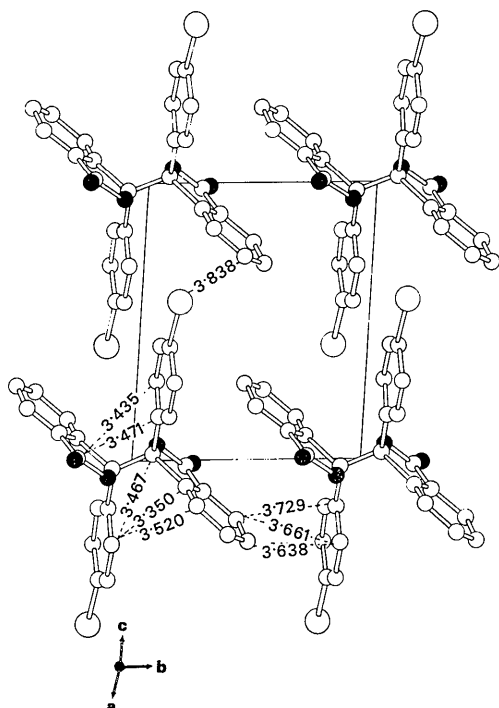


Fig. 5. View of the structure along [100] to illustrate the packing arrangement. Short intermolecular and intramolecular contacts are also indicated.

Table 7. Short non-bonded intramolecular distances

C(3)–C(14')	3.435 Å
C(4)–C(14')	3.471
C(5)–C(7')	3.467
C(5)–C(9')	3.520
C(5)–C(8')	3.350

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References

- ASHER, J. D. M. & SIM, G. A. (1965). *J. Chem. Soc.* p. 1584.
 BLOW, D. M. (1960). *Acta Cryst.* **13**, 168.
 BOND, W. L. (1959). *Acta Cryst.* **12**, 375.
 BUSING, W. R. & LEVY, H. A. (1964). *Acta Cryst.* **17**, 142.
 CROMER, D. T. & WABER, J. T. (1965). *Acta Cryst.* **18**, 104.
 CURL, R. F. (1959). *J. Chem. Phys.* **30**, 1529.
 DEWAR, M. J. S. & SCHMEISING, H. N. (1959). *Tetrahedron*, **5**, 166.
 FORSYTH, J. B. & WELLS, M. (1959). *Acta Cryst.* **12**, 412.
 FRIDRICHSONS, J. & MATHIESON, A. MCL. (1962). *Acta Cryst.* **15**, 119.
 FRIDRICHSONS, J., MATHIESON, A. MCL. & SUTOR, D. J. (1963). *Acta Cryst.* **16**, 1075.
 GABE, E. J. (1962). *Acta Cryst.* **15**, 759.
 HAMILTON, W. C. (1959). *Acta Cryst.* **12**, 609.
 HOWELLS, E. R., PHILLIPS, D. C. & ROGERS, D. (1950). *Acta Cryst.* **3**, 210.
 KANNAN, K. K., VIJAYAN, M. & NAMBU DIRI, M. M. G. (1966). *A Unified Set of Crystallographic Programmes*. Bangalore: Indian Institute of Science. To be published.
 KARLE, I. & KARLE, J. (1966). *Acta Cryst.* **20**, 555.
 KARTHA, G. & HAAS, D. J. (1964). *J. Amer. Chem. Soc.* **86**, 3630.
 LIPSON, H. & COCHRAN, W. (1957). *The Determination of Crystal Structures*. London: Bell.
 MATHIESON, A. MCL. & TAYLOR, J. C. (1963). *Acta Cryst.* **16**, 524.
 MCCONNELL, J. F., MATHIESON, A. MCL. & SCHOENBORN, B. P. (1964). *Acta Cryst.* **17**, 472.
 MILLS, O. S. & ROLLETT, J. S. (1961). In *Computing Methods and the Phase Problem in X-ray Crystal Analysis*, p. 117. Oxford: Pergamon Press.
 PRZYBYLSKA, M. & AHMED, F. R. (1958). *Acta Cryst.* **11**, 718.
 SCHOENBORN, B. P. & MCCONNELL, J. F. (1962). *Acta Cryst.* **15**, 779.
 SUTTON, L. E. (1958). *Tables of Interatomic Distances and Configuration in Molecules and Ions*. London: The Chemical Society.
 WILSON, A. J. C. (1942). *Nature, Lond.* **150**, 151.
 WILSON, A. J. C. (1949). *Acta Cryst.* **2**, 318.